

## 19. CONVOLUTION

**19.1. Superposition of infinitesimals: the convolution integral.**

The system response of an LTI system to a general signal can be reconstructed explicitly from the unit impulse response.

To see how this works, start with an LTI system represented by a linear differential operator  $L$  with constant coefficients. The system response to a signal  $f(t)$  is the solution to  $Lx = f(t)$ , subject to some specified initial conditions. To make things uniform it is common to specify “rest” initial conditions:  $x(t) = 0$  for  $t < 0$ .

We will approach this general problem by decomposing the signal into small packets. This means we partition time into intervals of length say  $\Delta t$ :  $t_0 = 0, t_1 = \Delta t, t_2 = 2\Delta t$ , and generally  $t_k = k\Delta t$ . The  $k$ th signal packet is the null signal (i.e. has value zero) except between  $t = t_k$  and  $t = t_{k+1}$ , where it coincides with  $f(t)$ . Write  $f_k(t)$  for the  $k$ th packet. Then  $f(t)$  is the sum of the  $f_k(t)$ 's.

Now by superposition the system response (with rest initial conditions) to  $f(t)$  is the sum of the system responses to the  $f_k(t)$ 's separately.

The next step is to estimate the system response to a single packet, say  $f_k(t)$ . Since  $f_k(t)$  is concentrated entirely in a small neighborhood of  $t_k$ , it is well approximated as a rate by a multiple of the delta function concentrated at  $t_k$ ,  $\delta(t - t_k)$ . The multiple should be chosen so that the cumulative totals match up; that is, it should be the integral under the graph of  $f_k(t)$ , which is itself well approximated by  $f(t_k)\Delta t$ . Thus we replace  $f_k(t)$  by

$$f(t_k) \cdot \Delta t \cdot \delta(t - t_k).$$

The system response to this signal, a multiple of a shift of the unit impulse, is the same multiple of the same shift of the weight function (= unit impulse response):

$$f(t_k) \cdot \Delta t \cdot w(t - t_k).$$

By superposition, adding up these packet responses over the packets which occur before the given time  $t$  gives the system response to the signal  $f(t)$  at time  $t$ . As  $\Delta t \rightarrow 0$  this sum approximates an integral taken over time between time zero and time  $t$ . Since the symbol  $t$  is already in use, we need to use a different symbol for the variable in the integral; let's use the Greek equivalent of  $t$ ,  $\tau$  (“tau”). The  $t_k$ 's get

replaced by  $\tau$  in the integral, and  $\Delta t$  by  $d\tau$ :

$$(1) \quad \boxed{x(t) = \int_0^t f(\tau)w(t - \tau) d\tau}$$

This is a really wonderful formula. Edwards and Penney call it “Duhamel’s principle,” but they seem somewhat isolated in this. Perhaps a better name would be the “superposition integral,” since it is no more and no less than an integral expression of the principle of superposition. It is commonly called the **convolution integral**. It describes the solution to a general LTI equation  $Lx = f(t)$  subject to rest initial conditions, in terms of the unit impulse response  $w(t)$ . Note that in evaluating this integral  $\tau$  is always less than  $t$ , so we never encounter the part of  $w(t)$  where it is zero.

**19.2. Example: the build up of a pollutant in a lake.** Every good formula deserves a particularly illuminating example, and perhaps the following will serve for the convolution integral. It is illustrated by the Mathlet **Convolution: Accumulation**. We have a lake, and a pollutant is being dumped into it, at a certain variable rate  $f(t)$ . This pollutant degrades over time, exponentially. If the lake begins at time zero with no pollutant, how much is in the lake at time  $t > 0$ ?

The exponential decay is described as follows. If a quantity  $p$  of pollutant is dropped into the lake at time  $\tau$ , then at a later time  $t$  it will have been reduced in amount to  $pe^{-a(t-\tau)}$ . The number  $a$  is the decay constant, and  $t - \tau$  is the time elapsed. We apply this formula to the small drip of pollutant added between time  $\tau$  and time  $\tau + \Delta\tau$ . The quantity is  $p = f(\tau)\Delta\tau$  (remember,  $f(t)$  is a *rate*; to get a *quantity* you must multiply by time), so at time  $t$  the this drip has been reduced to the quantity

$$e^{-a(t-\tau)} f(\tau)\Delta\tau$$

(assuming  $t > \tau$ ; if  $t < \tau$ , this particular drip contributed zero). Now we add them up, starting at the initial time  $\tau = 0$ , and get the convolution integral (1), which here is

$$(2) \quad x(t) = \int_0^t f(\tau)e^{-a(t-\tau)} d\tau.$$

We found our way straight to the convolution integral, without ever mentioning differential equations. But we can also solve this problem by setting up a differential equation for  $x(t)$ . The amount of this chemical in the lake at time  $t + \Delta t$  is the amount at time  $t$ , minus the fraction

that decayed, plus the amount newly added:

$$x(t + \Delta t) = x(t) - ax(t)\Delta t + f(t)\Delta t$$

Forming the limit as  $\Delta t \rightarrow 0$ , we obtain

$$(3) \quad \dot{x} + ax = f(t), \quad x(0) = 0.$$

We conclude that (2) gives us the solution with rest initial conditions.

An interesting case occurs if  $a = 0$ . Then the pollutant doesn't decay at all, and so it just builds up in the lake. At time  $t$  the total amount in the lake is just the total amount dumped in up to that time, namely

$$\int_0^t f(\tau) d\tau,$$

which is consistent with (2).

**19.3. Convolution as a “product”.** The integral (1) is called the *convolution* of  $w(t)$  and  $f(t)$ , and written using an asterisk:

$$(4) \quad w(t) * f(t) = \int_0^t w(t - \tau)f(\tau) d\tau, \quad t > 0.$$

We have now fulfilled the promise we made at the beginning of Section 18: we can explicitly describe the system response, with rest initial conditions, to any input signal, if we know the system response to just one input signal, the unit impulse:

**Theorem.** The solution to an LTI equation  $Lx = f(t)$ , of any order, with rest initial conditions, is given by

$$x(t) = w(t) * f(t),$$

where  $w(t)$  is the unit impulse response.

If  $L$  is an LTI differential operator, we should thus be able to reconstruct its characteristic polynomial  $p(s)$  (so that  $L = p(D)$ ) from its unit impulse response. This is one of the things the Laplace transform does for us; in fact, the Laplace transform of  $w(t)$  is the reciprocal of  $p(s)$ : see Section 21.

The expression (4) can be interpreted formally by a process known as “flip and drag.” It is illustrated in the Mathlet **Convolution: Flip and Drag**.

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